# Computational Studies on the Reaction Pathways of $\mathrm{CF}_{3} \mathrm{Br}$ with $\mathbf{O}\left({ }^{1} \mathbf{D},{ }^{\mathbf{3}} \mathbf{P}\right)$ Atoms 

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#### Abstract

The triplet and singlet potential energy surfaces (PES) for $\mathrm{O}\left({ }^{3} \mathrm{P},{ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3} \mathrm{Br}$ reactions have been studied using the density functional method at the B3LYP level. Geometries, energies, and vibrational frequencies of $\mathrm{CF}_{3} \mathrm{Br}, \mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right.$ and $\left.{ }^{3} \mathrm{~A}^{\prime \prime}\right), \mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right), \mathrm{CF}_{3}, \mathrm{BrO}, \mathrm{CF}_{2} \mathrm{O}, \mathrm{BrF}$, and $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ complexes have been examined. Transition states connecting these species have been characterized and the whole potential surfaces could satisfactorily describe the $\mathrm{O}+\mathrm{CF}_{3} \mathrm{Br}$ reactions. In the gas phase, the OBr radical is the main product on both the singlet and triplet PESs. In a condensed matrix, however, $\mathrm{CF}_{3} \mathrm{OBr}$ should be the dominant product generating from the recombination of nascent $\mathrm{CF}_{3}+\mathrm{OBr}$.


## 1. Introduction

$\mathrm{CF}_{3} \mathrm{Br}$ (Halon 1301) has long been used as fire fighting agent and refrigerant, and is believed to contribute significantly to atmospheric bromine. ${ }^{1}$ Since $\mathrm{CF}_{3} \mathrm{Br}$ does not react with OH radicals, its tropospheric lifetimes is very long (65-77 years). ${ }^{2}$ In the stratosphere, $\mathrm{CF}_{3} \mathrm{Br}$ can either be photodecomposed or react directly with atomic oxygen to generate ozone-depleting agents such as Br and BrO in the gas phase. The following reactions between oxygen atoms and $\mathrm{CF}_{3} \mathrm{Br}$ have been extensively investigated in the gas phase: ${ }^{3-5}$

$$
\begin{gather*}
\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3} \mathrm{Br} \rightarrow \mathrm{CF}_{3}+\mathrm{BrO}  \tag{1}\\
\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3} \mathrm{Br} \rightarrow \mathrm{Br}+\mathrm{CF}_{3} \mathrm{O}  \tag{2}\\
\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3} \mathrm{Br} \rightarrow \mathrm{CF}_{3} \mathrm{Br}+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \tag{3}
\end{gather*}
$$

Thompson and Ravishankara ${ }^{3}$ have studied the branching ratio for the reactive loss and quenching of $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ by monitoring the time-resolved VUV atomic resonance fluorescence of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$. Recently, Cronkhite and Wine ${ }^{4}$ pointed out that reaction 1 is the dominant one for the competing reactions 1 and 2. Alagia et al. ${ }^{5}$ have built a model potential energy surface coupled with their molecular beam scattering experiments for $\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3}$ Br system. They suggested that BrO in reaction 1 could be formed on a reactive PES with a deep energy well stabilizing $\mathrm{CF}_{3} \mathrm{OBr}$ intermediates. However, $\mathrm{CF}_{3} \mathrm{OBr}$ has never been observed in the gas phase.

Lorenzen-Schmidt et al. ${ }^{6}$ and Minkwitz et al. ${ }^{7}$ have successfully isolated $\mathrm{CF}_{3} \mathrm{OBr}$ and recorded its IR and Raman spectra in an Ar matrix as well as in pure solid state. Quite recently, our group reinvestigated the IR spectra of $\mathrm{CF}_{3} \mathrm{OBr}$ in an Ar matrix, ${ }^{8}$ and infrared absorptions at $1249.4,1204.2$, and 1201.9 $\mathrm{cm}^{-1}$ were assigned to $\mathrm{CF}_{3} \mathrm{OBr}$ with the aid of density functional calculations at the B3LYP/6-31+G(d) level. This work is now completed by detailed quantum chemical MO calculations of the potential energy surface (PES) for reactions between $\mathrm{CF}_{3}{ }^{-}$ Br and oxygen (both ${ }^{1} \mathrm{D}$ and ${ }^{3} \mathrm{P}$ states). Geometries, energies, vibrational frequencies of reactants, products, reaction intermediates, and transition states (TS) are obtained, and we hope

[^0]these results will be valuable for better understanding the reaction mechanism of the $\mathrm{CF}_{3} \mathrm{Br}+\mathrm{O}$ system.

## 2. Computation Methods

Quantum chemical MO calculations were carried out using Gaussian 98 codes. ${ }^{9}$ Initial potential energy surface scan was performed using the B3LYP ${ }^{10,11}$ method with the $6-31+G(d)$ basis set. Stationary points of interests were further optimized using the $6-311+\mathrm{G}(\mathrm{d})$ basis set. Harmonic vibrational frequencies were calculated using analytical second derivatives, and zero point vibrational energies (ZPE) were derived. The number of imaginary frequencies were used to characterize the nature of located structures. The ZPEs and the vibrational frequencies of reaction intermediates and transition states are used without scaling. Single point energy calculations were done with $6-311+G(2 d f)$ basis sets using the $6-311+G(d)$ geometries.

Transition state optimizations were done with the synchronous transit-guided quasi-Newton (STQN) method ${ }^{12}$ at the B3LYP/ $6-311+G(d)$ level. The program can generate a guess for the transition structure that is midway between the reactant and product, and it then goes on to optimize that starting structure to a first-order saddle point automatically. At all the calculated transition states, vibrational normal modes were analyzed to testify the nature of the saddle point. In some cases, intrinsic reaction coordinate (IRC) calculations were performed to follow the reaction path.

## 3. Results and Discussion

In section 3.1, we will first discuss the intersystem crossing between the singlet and triplet potential energy surfaces (PESs) of the $\mathrm{O}+\mathrm{CF}_{3} \mathrm{Br}$ system; then we will describe the PESs in detail. These profiles are shown in Figures 1 and 2. Reaction intermediates and the transition states (TS) are calculated with their energies in Table 1 and geometries in Figures 3 and 4. Finally, we will discuss the vibrational frequencies (Tables 2 and 3) of the reaction intermediates and TS in section 3.2.
3.1. Potential Energy Surfaces. Energies of $O\left({ }^{3} \mathrm{P}\right)$ and $O$ $\left.{ }^{l} D\right)$. Before we proceed to discuss the reaction PESs, it is necessary to confirm the veracity of our computation methods. In Table 1, experimental energies ${ }^{13,14}$ of some species are included for comparison. We note the energy difference between $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ and $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ is calculated to be $63.4 \mathrm{kcal} / \mathrm{mol}$, which severely deviates from the measured value of $45.4 \mathrm{kcal} / \mathrm{mol} .{ }^{13}$ This large difference originates from the wrong symmetry of the wave function for the $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ atom in a simple closed-shell


Figure 1. Relative energy of $\mathrm{O}+\mathrm{CF}_{3} \mathrm{Br}$ interaction on the singlet and triplet PESs plotted as a function of $\mathrm{Br}-\mathrm{O}$ separation with bent configuration. Calculated at the B3LYP/6-31+G(d) level.


Figure 2. Singlet (solid lines) and triplet (dotted lines) potential energy profiles of the $\mathrm{CF}_{3} \mathrm{Br}+\mathrm{O}$ system, plotted using B3LYP/6-311+G(d) energies plus ZPE corrections.

TABLE 1: Relative Energies (kcal/mol) of the Reaction Intermediates and Transition States

|  | energies |  |  |  |
| :--- | :---: | :---: | :---: | :---: |
| species | $6-311+\mathrm{G}(\mathrm{d})$ | G(d) +ZPE <br>  <br>  <br> $\mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ <br> $\mathrm{G}(2 \mathrm{df})^{a}$ | expt $^{b}$ |  |
| $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ (end on) | -71.36 | -68.80 | -75.24 |  |
| $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})$ (side on) | -66.34 | -64.88 | -69.31 |  |
| $\mathrm{CF}_{2} \mathrm{O}+\mathrm{BrF}$ | -64.28 | -63.08 | -67.18 |  |
| $\mathrm{CF}_{3} \mathrm{O}+\mathrm{Br}$ | -62.88 | -62.06 | -66.07 |  |
| $\mathrm{CF}_{2} \mathrm{O}+\mathrm{Br}+\mathrm{F}$ | -28.75 | -27.78 | -28.99 | -34.2 |
| $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ | -6.84 | -6.95 | -5.26 | -10.2 |
| $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ | -6.45 | -5.30 | -14.04 |  |
| $\mathrm{CF}_{3} \mathrm{Br}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | -1.35 | -1.13 | -1.55 |  |
| $\mathrm{CF}_{3}+\mathrm{BrO}$ | 0 | 0 | 0 | 0 |
| $\mathrm{CF}_{3} \mathrm{Br}+\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ | 12.53 | 12.13 | 8.05 | 14.0 |
| TS 1 | 63.36 | 63.36 | 63.41 | 45.45 |
| TS 2 | 25.86 | 25.68 | 22.96 |  |
| TS 3 | 25.35 | 26.18 | 17.47 |  |
|  | -16.20 | -15.28 | -17.34 |  |

${ }^{a}$ Single point energy calculated at the B3LYP/6-311+G(d) geometries ${ }^{b}$ Relative values of heats of formation $\Delta H_{\mathrm{f}}$ (298K). The data are from ref 13 , except that of $\mathrm{CF}_{3} \mathrm{O}$, which is from ref 14 .
type Kohn-Sham functional. This could be fixed by introducing perfect-pairing general valence bond (GVB-PP) calculation, ${ }^{9}$ but unfortunately, it is not available with B3LYP DFT methods.

The singlet-triplet intersystem crossing we are going to discuss is only a qualitative feature, and the aforementioned incorrect energy gap between $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ and $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ will not influence the results. On the other hand, the calculated relative energies accord with experiments with a uncertainty of $2-5 \mathrm{kcal} / \mathrm{mol}$, as shown in Table 1. Therefore, higher level calculations such as $\operatorname{CCSD}(\mathrm{t})$ or QCISD have not been performed to correct the energies.


Figure 3. Geometries of $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime},{ }^{3} \mathrm{~A}^{\prime \prime}\right), \mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right), \mathrm{CF}_{2} \mathrm{O}-$ ( BrF ) side-on complex, and $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ end-on complex. The intermediates all pertain to $C_{s}$ symmetry.


Figure 4. Calculated geometries of the reaction transition states. The TSs all have $C_{s}$ symmetry.

TABLE 2: Calculated Vibrational Frequencies ( $\mathrm{cm}^{-1}$ ), IR Intensities (in Parentheses, $\mathrm{km} / \mathrm{mol}$ ) of $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime},{ }^{3} \mathrm{~A}^{\prime \prime}\right.$ ), $\mathrm{CF}_{3} \mathrm{OBr}$, and $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ Complexes at the B3LYP/ 6-311+G(d) Level

| $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ | $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ | $\mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ | $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})($ side $)$ | $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}(\mathrm{end})$ |
| :---: | :---: | :---: | :---: | :---: |
| $1224.1(317)$ | $1188.9(312)$ | $1214.0(327)$ | $1970.7(485)$ | $1937.1(772)$ |
| $1210.2(291)$ | $1180.1(278)$ | $1185.1(667)$ | $1205.5(455)$ | $1245.7(447)$ |
| $1043.7(482)$ | $1047.0(597)$ | $1163.3(386)$ | $956.8(61)$ | $975.0(66)$ |
| $735.9(50)$ | $750.2(57)$ | $899.2(6)$ | $768.2(58)$ | $775.9(40)$ |
| $696.0(43)$ | $541.7(1)$ | $719.2(1)$ | $642.8(35)$ | $627.1(64)$ |
| $542.9(0)$ | $541.5(1)$ | $644.9(3)$ | $615.4(6)$ | $624.6(11)$ |
| $525.6(2)$ | $332.4(1)$ | $599.8(3)$ | $576.3(5)$ | $580.7(5)$ |
| $318.6(3)$ | $297.3(0)$ | $518.0(1)$ | $90.7(2)$ | $124.6(3)$ |
| $262.3(0)$ | $296.4(0)$ | $423.7(0)$ | $63.4(2)$ | $120.1(2)$ |
| $249.1(1)$ | $104.5(7)$ | $337.3(3)$ | $50.3(1)$ | $90.9(6)$ |
| $138.8(8)$ | $27.7(0)$ | $169.5(1)$ | $36.9(1)$ | $47.6(0)$ |
| $31.8(9)$ | $15.7(0)$ | $88.0(1)$ | $34.1(1)$ | $37.8(0)$ |

TABLE 3: Calculated Vibrational Frequencies ( $\mathrm{cm}^{-1}$ ) of the Reaction Transition States at the B3LYP/6-311 $+\mathbf{G}(\mathbf{d})$ Level

| TS1 |  |  |
| :---: | :---: | :---: |
| $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right) \rightarrow$ <br> $\mathrm{CF}_{3}+\mathrm{BrO}$ | TS 2 <br> $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right) \rightarrow$ <br> $\mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ | TS 3 <br> $\mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right) \rightarrow$ <br> $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})($ side on $)$ |
| 480.3 i | 437.1 i | 757.8 i |
| 63.7 | 94.3 | 34.9 |
| 150.6 | 261.6 | 239.5 |
| 168.1 | 301.2 | 284.8 |
| 241.8 | 332.2 | 316.5 |
| 287.3 | 536.3 | 479.4 |
| 518.3 | 563.5 | 586.0 |
| 518.8 | 585.1 | 604.1 |
| 679.4 | 708.5 | 705.6 |
| 972.5 | 904.2 | 941.9 |
| 1218.2 | 1139.8 | 1193.1 |
| 1226.2 | 1326.2 | 1429.0 |

Intersystem Crossing. To characterize the seam of crossing, we scanned the PES for interaction between $\mathrm{O}\left({ }^{1} \mathrm{D},{ }^{3} \mathrm{P}\right)$ and $\mathrm{CF}_{3}{ }^{-}$ Br at the B3LYP/6-31+G(d) level. Previous calculations ${ }^{15,16}$ showed that O approaching $\mathrm{CF}_{3} \mathrm{I}$ and $\mathrm{CH}_{3} \mathrm{I}$ in a bent configu-
ration is energetically favored. Similarly, both singlet and triplet $\mathrm{CF}_{3} \mathrm{BrO}$ complexes have bent geometries; therefore we have only calculated the energy along the bent collision trajectories.

As shown in Figure 1, energies along the interaction path are plotted as a function of $\mathrm{CF}_{3} \mathrm{Br}-\mathrm{O}$ separation. During optimization, we kept the $\mathrm{CF}_{3} \mathrm{Br}$ part unchanged, and varied the $\mathrm{Br}-\mathrm{O}$ bond length. The O atom approaches the Br side in $\mathrm{CF}_{3} \mathrm{Br}$ with a $\mathrm{C}-\mathrm{Br}-\mathrm{O}$ angle of $106.8^{\circ}$, the same angle as in the $\mathrm{CF}_{3} \mathrm{BrO}$ adducts. As seen from Figure 1, the ISC occurs at a $\mathrm{Br}-\mathrm{O}$ distance of about $2.1 \AA$, which is midway between $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ and $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$. The crossing point is $2.7 \mathrm{kcal} /$ mol above the $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{CF}_{3} \mathrm{Br}$ reactants, and also 5.1 and 12.6 $\mathrm{kcal} / \mathrm{mol}$ higher than the $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ and $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$, respectively. The energy of $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ depends significantly on the basis set, as can be seen in Table 1. This implies that the ISC point would also move when using a better basis set. Since the level of theory employed here is below that used to calculate the whole PESs, this picture for ISC is expected to be qualitatively correct.

Previous calculations on the reactions between $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ and $\mathrm{CF}_{3} \mathrm{I},{ }^{15} \mathrm{CH}_{3} \mathrm{I},{ }^{16}$ and $\mathrm{C}_{2} \mathrm{H}_{5} \mathrm{I}^{17}$ have shown that the singlet-triplet PES intersections are also important features. As far as we know, the ISC for the $\mathrm{CF}_{3} \mathrm{Br}+\mathrm{O}$ system is similar to that of the $\mathrm{C}_{2} \mathrm{H}_{5} \mathrm{I}$ $+\mathrm{O}\left({ }^{1} \mathrm{D},{ }^{3} \mathrm{P}\right)$ system, ${ }^{17}$ of which the ISC occurs at the exit of the entrance valley on the PES.
$\mathrm{CF}_{3} \mathrm{BrO}$ and $\mathrm{CF}_{3} \mathrm{OBr}$. The overall potential energy surfaces are given in Figure 2. Starting from $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ and $\mathrm{CF}_{3} \mathrm{Br}$, a vdW complex $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ is located with a binding energy of 1.1 $\mathrm{kcal} / \mathrm{mol}(\mathrm{B} 3 \mathrm{LYP} / 6-311+\mathrm{G}(\mathrm{d})+$ ZPE). Interaction between $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ and $\mathrm{CF}_{3} \mathrm{Br}$ leads directly to $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$, which is stabilized by about $68 \mathrm{kcal} / \mathrm{mol}$ with respect to $\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3^{-}}$ Br . Singlet $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ is a true covalently bound molecule, and it is $5.1 \mathrm{kcal} / \mathrm{mol}$ lower in energy than $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$. As can be seen from Figure 3, the $\mathrm{Br}-\mathrm{O}$ bond length in $\mathrm{CF}_{3} \mathrm{BrO}$ $\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ is $2.776 \AA$, indicating a weak interaction. When it comes to $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$, the $\mathrm{Br}-\mathrm{O}$ bond length becomes $1.723 \AA$, which is even shorter than that in $\mathrm{OBr}(1.756 \AA)$, and the $\mathrm{C}-\mathrm{Br}$ is enlarged to $2.051 \AA$ from $1.948 \AA$ in free $\mathrm{CF}_{3} \mathrm{Br}$. Mulliken population reveals that the BrO bond in $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ is ionic in nature with +0.578 atomic charge on Br and -0.497 on O . This may arise from the donation of the bromine lone pair electrons to the vacant p orbital on oxygen.

Dissociation of $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ to doublet $\mathrm{CF}_{3}+\mathrm{BrO}$ proceeds via TS1, a transition state lying $24.5 \mathrm{kcal} / \mathrm{mol}$ above the minimum $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ and $13.6 \mathrm{kcal} / \mathrm{mol}$ above $\mathrm{CF}_{3}+\mathrm{BrO}$. On the singlet surface, $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ generated from $\mathrm{O}\left({ }^{1} \mathrm{D}\right)+$ $\mathrm{CF}_{3} \mathrm{Br}$ is initially excited by nearly $70 \mathrm{kcal} / \mathrm{mol}$ in energy. This molecule can undergo further dissociation/isomerization pathways on the PES, as shown in Figure 2. First, dissociation of $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ by $\mathrm{C}-\mathrm{Br}$ bond fission to form $\mathrm{CF}_{3}+\mathrm{OBr}$ has no distinct barrier, and the process needs only $18 \mathrm{kcal} / \mathrm{mol}$ energy input. Previous experiment ${ }^{4}$ found that BrO formation via reaction 1 is the dominant channel with a branching ratio of 0.49 at 298 K . This result accords with our PES qualitatively. Second, $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ could also isomerize through TS2 to form $\mathrm{CF}_{3} \mathrm{OBr}$. This three-center TS lies $31.5 \mathrm{kcal} / \mathrm{mol}$ above $\mathrm{CF}_{3}-$ $\mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$. However, in comparison with the dissociation channel leading to $\mathrm{CF}_{3}+\mathrm{BrO}$, this isomerization process is less energetically favored. The geometric parameters of TS1 and TS2 are given in Figure 4.

At the present stage, it is of interest to discuss the origin of $\mathrm{CF}_{3} \mathrm{OBr}$ in the matrix isolation experiment. ${ }^{6,8}$ Lorenzen-Schmidt et al. ${ }^{6}$ have addressed that $\mathrm{CF}_{3} \mathrm{OBr}$ came from photodecomposition of the $\mathrm{CF}_{3} \mathrm{Br}-\mathrm{O}_{3}$ complex, and the process is predominantly an insertion reaction. They suggested that direct insertion mechanism should apply or else the byproduct like $\mathrm{O}_{2}$ might compete with $\mathrm{CF}_{3}$ and BrO to give $\mathrm{CF}_{3} \mathrm{OO}$ or other products. However, under our matrix isolation experiment conditions, ${ }^{8}$ the reaction precursors $\left(\mathrm{CF}_{3} \mathrm{Br}\right.$ and $\left.\mathrm{O}_{3}\right)$ were in good isolation and the product $\mathrm{CF}_{3} \mathrm{OBr}$ resulted from separate $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ and $\mathrm{CF}_{3}$ Br pairs, so the competing pathways involving $\mathrm{O}_{2}$ can be totally excluded. Accordingly, it is reasonable to recognize that interaction between $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ and $\mathrm{CF}_{3} \mathrm{Br}$ leads to $\mathrm{CF}_{3} \mathrm{BrO}$ adduct, which readily dissociates to $\mathrm{CF}_{3}$ and BrO , and fast recombination of nascent $\mathrm{CF}_{3}$ and BrO leads to $\mathrm{CF}_{3} \mathrm{OBr}$ formation.

A dissociation/recombination mechanism for $\mathrm{O}_{3}$ and $\mathrm{CF}_{3} \mathrm{Br}$ reactions in matrix can be proposed as follows:

$$
\begin{gathered}
\mathrm{O}_{3}+h v \rightarrow \mathrm{O} *\left({ }^{1} \mathrm{D}\right)+\mathrm{O}_{2}\left({ }^{1} \Delta_{\mathrm{g}}\right) \text { (in cage) } \\
\mathrm{O} *\left({ }^{1} \mathrm{D}\right) \rightarrow \mathrm{O}\left({ }^{1} \mathrm{D}\right) \text { (cage escape and migration) } \\
\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3} \mathrm{Br} \rightarrow \mathrm{CF}_{3} \mathrm{BrO} \text { (addition) } \\
\mathrm{CF}_{3} \mathrm{BrO} \rightarrow \mathrm{CF}_{3}+\mathrm{OBr} \text { (dissociation, fast) } \\
\mathrm{CF}_{3}+\mathrm{OBr} \rightarrow \mathrm{CF}_{3} \mathrm{OBr} \text { (recombination, fast) }
\end{gathered}
$$

This mechanism can satisfactorily explain the reaction products observed in matrix IR spectra. Moreover, it also supports the gas-phase experiments. ${ }^{3-5}$

Final Product: $\mathrm{CF}_{2} \mathrm{O}$ Formation. Now we will concentrate on the latter region of the PES where the final products $\mathrm{CF}_{2} \mathrm{O}$ and BrF were produced from $\mathrm{O}\left({ }^{1} \mathrm{D}\right)+\mathrm{CF}_{3} \mathrm{Br}{ }^{6,8}$ reactions. This channel is important in condensed phase reaction, ${ }^{6,8}$ although it is not observed in gas phase.

Obviously, $\mathrm{CF}_{2} \mathrm{O}$ should come directly from $\mathrm{CF}_{3} \mathrm{OBr}$ or $\mathrm{CF}_{3} \mathrm{O}$ in which O binds directly with C . We will consider $\mathrm{CF}_{3} \mathrm{O}$ first. About $41 \mathrm{kcal} / \mathrm{mol}$ energy input is required to break the $\mathrm{O}-\mathrm{Br}$ bond in $\mathrm{CF}_{3} \mathrm{OBr}$ to form $\mathrm{CF}_{3} \mathrm{O}+\mathrm{Br}$. Since minor amounts of $\mathrm{CF}_{2} \mathrm{O}+\mathrm{BrF}$ were observed in matrix isolation experiments, ${ }^{6,8}$ it seems a small fraction of initially excited $\mathrm{CF}_{3} \mathrm{OBr}$ could still dissociate to $\mathrm{CF}_{3} \mathrm{O}+\mathrm{Br}$. One possible channel is that $\mathrm{CF}_{2} \mathrm{O}+$ BrF could originate from the following steps: dissociation of $\mathrm{CF}_{3} \mathrm{O}$ to $\mathrm{CF}_{2} \mathrm{O}+\mathrm{F}$, followed by F atom recombination with Br in a matrix cage forming BrF finally. Normally, $\mathrm{C}-\mathrm{F}$ bonds are very strong, but it is interesting to notice that the third $\mathrm{C}-\mathrm{F}$ bond in $\mathrm{CF}_{3} \mathrm{O}$ is much weaker than the other two. B3LYP/6$311+\mathrm{G}(\mathrm{d})$ calculations predict its bond energy to be only 21 $\mathrm{kcal} / \mathrm{mol}$, much lower than the normal $\mathrm{C}-\mathrm{F}$ bond strength of $110-120 \mathrm{kcal} / \mathrm{mol}$.

Another possible channel is Br abstraction of F in $\mathrm{CF}_{3} \mathrm{O}$ to form a side-on $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})$ complex. We have explored the PES and found a transition state (TS3) for this Br abstraction channel:

$$
\mathrm{Br}+\mathrm{CF}_{3} \mathrm{O} \rightarrow \mathrm{TS} 3 \rightarrow \mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})
$$

Figure 4 gives the structure of TS3. In this transition state, the F atom (abstracted by Br ) is $1.649 \AA$ away from the C atom, which is $0.32 \AA$ longer than normal $\mathrm{C}-\mathrm{F}$ bond. The $\mathrm{Br}-\mathrm{F}$ distance is $2.154 \AA$, still $0.3 \AA$ away from the free BrF molecule ( $1.809 \AA$ ). TS3 is $12.5 \mathrm{kcal} / \mathrm{mol}$ above $\mathrm{CF}_{3} \mathrm{O}+\mathrm{Br}$. The final product of this channel is a side-on complex, $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})$, which is $1.0 \mathrm{kcal} / \mathrm{mol}$ more stable than separated $\mathrm{CF}_{2} \mathrm{O}+\mathrm{BrF}$. Geometries of $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})$ complexes can be found in Figure 3.

The unimolecular dissociation of $\mathrm{CF}_{3} \mathrm{OBr}$ to form the $\mathrm{CF}_{2} \mathrm{O}-$ BrF complex is another reaction pathway generating $\mathrm{CF}_{2} \mathrm{O}$. Breaking the $\mathrm{C}-\mathrm{F}$ and $\mathrm{O}-\mathrm{Br}$ bonds in $\mathrm{CF}_{3} \mathrm{OBr}$ requires lots of energy input. The $\mathrm{C}-\mathrm{F}$ bond in $\mathrm{CF}_{3} \mathrm{OBr}$ should not be as weak as that in $\mathrm{CF}_{3} \mathrm{O}$ (vide supra), although it is hard to determine the $\mathrm{C}-\mathrm{F}$ bond strength here. On that basis, a TS which looks like $\mathrm{CF}_{2} \mathrm{OBr}$ with a separate F atom attacking at Br must be high in energy on the PES. Several attempts trying to locate this TS connecting $\mathrm{CF}_{3} \mathrm{OBr}$ and $\mathrm{CF}_{2} \mathrm{O}(\mathrm{BrF})$ (side-on complex) eventually failed. Thus, we could not judge for the present time whether this unimolecular reaction is applicable, and whether there is a TS for this reaction.

There are two conformers of $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ complexes, namely, side-on and end-on, whose structures are shown in Figure 3. The side-on complex is $1 \mathrm{kcal} / \mathrm{mol}$ more stable than separated $\mathrm{CF}_{2} \mathrm{O}+\mathrm{BrF}$, and the $\mathrm{Br}-\mathrm{F}$ unit in this complex is parallel to the $\mathrm{CF}_{2} \mathrm{O}$ planar with Br interacting with O atom and F with central C atom. The end-on complex has a planar structure with terminal O atom linking $\mathrm{Br}-\mathrm{F}$ through the bromine atom. The end-on complex is $3 \mathrm{kcal} / \mathrm{mol}$ lower in energy than isolated $\mathrm{CF}_{2} \mathrm{O}+\mathrm{BrF}$. Both complexes have been observed in matrix isolation experiment. ${ }^{6}$
3.2. Vibrational Frequencies. Minkwitz et al. have reported the results of MP2/6-31G* and HF/6-311G* calculations on $\mathrm{CF}_{3}{ }^{-}$ OBr. ${ }^{7}$ Recently, Kwon et al. ${ }^{18}$ calculated the structure, rotational barrier and vibrational frequencies of $\mathrm{CF}_{3} \mathrm{OBr}$ using B3LYP method. In our previous study on the IR frequencies of $\mathrm{CF}_{3}-$ $\mathrm{OBr},{ }^{8}$ we have already carried out B3LYP/6-31+G(d) calculations on both $\mathrm{CF}_{3} \mathrm{OBr}$ and $\mathrm{CF}_{3} \mathrm{BrO}$. In this paper, we have completed the previous study with higher level calculations on $\mathrm{CF}_{3} \mathrm{OBr}$ and $\mathrm{CF}_{3} \mathrm{BrO}$ (both singlet and triplet states). Moreover, final products such as $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ complexes and transition states were also calculated. Vibrational frequencies of these reaction intermediates and TSs are compiled in Tables 2 and 3.

For singlet $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$, the most intense bands are two $\mathrm{CF}_{3}$ asymmetric stretch (1224.1, $1210.2 \mathrm{~cm}^{-1}$ ) and one $\mathrm{CF}_{3}$ symmetric stretch $\left(1043.7 \mathrm{~cm}^{-1}\right)$. These modes are close to the corresponding vibrations in $\mathrm{CF}_{3}$ radical at 1223.0, 1222.9, and $1062.5 \mathrm{~cm}^{-1}$, suggesting a weak $\mathrm{C}-\mathrm{Br}$ bond in $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$. The B3LYP/6-311 $+\mathrm{G}(\mathrm{d})$ results show that the $\mathrm{C}-\mathrm{Br}$ bond in $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ is only $18 \mathrm{kcal} / \mathrm{mol}$ in strength, much lower than $70.6 \mathrm{kcal} / \mathrm{mol}$ in $\mathrm{CF}_{3} \mathrm{Br} .{ }^{19}$ For triplet $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$, the three $\mathrm{CF}_{3}$ stretch modes are also the most intense absorptions. However, these frequencies are very close to the corresponding modes in free $\mathrm{CF}_{3} \mathrm{Br}$ which are 1055.3, 1174.4, and 1174.4 $\mathrm{cm}^{-1}$. This implies that the interaction between $\mathrm{CF}_{3} \mathrm{Br}$ and O in $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ is quite weak.

Calculated vibrational frequencies of $\mathrm{CF}_{3} \mathrm{OBr}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ can reproduce our measured matrix isolation IR spectra ${ }^{8}$ very well. B3LYP/6-311+G(d) calculations give three $\mathrm{CF}_{3}$ stretch frequencies at $1214.0,1185.1$, and $1163.3 \mathrm{~cm}^{-1}$, slightly lower than the experimental values of $1249.4,1204.2$, and $1201.9 \mathrm{~cm}^{-1}$. For the final products of $\mathrm{CF}_{2} \mathrm{O}$ and its complexes with BrF , we also found a good match between theoretical and experimental results. Calculated $\mathrm{C}-\mathrm{O}$ stretch frequencies in $\mathrm{CF}_{2} \mathrm{O}, \mathrm{CF}_{2} \mathrm{O}-$ ( BrF ) (side on), and $\mathrm{CF}_{2} \mathrm{O}-\mathrm{BrF}$ (end on) molecules are at 1974.3, 1970.7, and $1937.1 \mathrm{~cm}^{-1}$, respectively. The values and the trend are also in good agreement with measured bands at 1913.3, 1912.5, and $1887.3 \mathrm{~cm}^{-1} .{ }^{6}$

## 4. Conclusion

Singlet and triplet potential energy surfaces for $\mathrm{O}\left({ }^{3} \mathrm{P},{ }^{1} \mathrm{D}\right)+$ $\mathrm{CF}_{3} \mathrm{Br}$ reaction pathways have been thoroughly explored by DFT method at B3LYP/6-311+G(d) level. Our calculations show that
interaction of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with $\mathrm{CF}_{3} \mathrm{Br}$ results in $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ formation directly, followed by dissociation to $\mathrm{CF}_{3}+\mathrm{BrO}$ via TS1. On the singlet PES, $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ and $\mathrm{CF}_{3} \mathrm{Br}$ interaction leads to $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$, and subsequent dissociation of this intermediate to form $\mathrm{CF}_{3}+\mathrm{BrO}$ is barrierless. Singlet-triplet intersystem crossing (ISC) is midway between $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{3} \mathrm{~A}^{\prime \prime}\right)$ and $\mathrm{CF}_{3} \mathrm{BrO}$ $\left({ }^{1} \mathrm{~A}^{\prime}\right)$ at the entrance channel on the PESs. The singlet adduct $\mathrm{CF}_{3} \mathrm{BrO}\left({ }^{1} \mathrm{~A}^{\prime}\right)$ could further isomerize via TS2 to form $\mathrm{CF}_{3}{ }^{-}$ OBr . Initially excited $\mathrm{CF}_{3} \mathrm{OBr}$ easily dissociates to $\mathrm{CF}_{3} \mathrm{O}+\mathrm{Br}$. Since the third $\mathrm{C}-\mathrm{F}$ bond in $\mathrm{CF}_{3} \mathrm{O}$ is calculated to be very weak, we believe either fission of the $\mathrm{C}-\mathrm{F}$ bond or Br attacking on the F atom (via TS3) would both generate $\mathrm{CF}_{2} \mathrm{O}$, a final product observed in matrix isolation experiments. The PESs obtained in this study can qualitatively explain the complicated behavior of $\mathrm{O}+\mathrm{CF}_{3} \mathrm{Br}$ reactions.

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Supporting Information Available: Tables of atomic coordinates. This material is available free of charge via the Internet at http://pubs.acs.org.

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